

*Schrödinger's Operator
and the
zeros of the eta function*

by

Mohammed Fulano

B.Sc. in Applied Mathematics

Summary

*Physicists study quantum systems, like free particles, using **Schrödinger** operator. Here, I will expose some ideas turning around the known relation between the roots of the zeta function and random matrices and their eigenvalues.*

3rd August, 2021

The **Schrödinger** equation is a linear operator, one of its parts acts as a **Hamiltonian** \hat{H} on a wavefunction Ψ , and is presented as a **Hermetian** matrix

$$\overline{H^T} = H$$

the wave function Ψ , gives essential information, that is measurements about a particle, these measurements (position, momentum, energy...) are known to be the eigenvalues of H . The squared amplitude of the wave function for $X = (x, y, z)$ and time t :

$$\|\Psi(x, y, z, t)\|^2$$

is the probability that a particle has a measurement M at a time t .

Ordinary Wave equation in two space dimensions has the form

$$\Delta u = \frac{1}{\alpha^2} u_{tt}$$

it has some relation with two other known equations namely :

$$\Delta u = 0$$

$$\Delta u = |k|u_t$$

Since there is an interesting relation between the roots of zeta function $\sigma_n + i\gamma_n$ and the eigenvalues of random matrices, a relation that is much studied by mathematicians and physicists since **Montgomery's** discovery, I try here , through it, to speculate on some possible relation between: the Riemann hypothesis **RH** , the potential V , and the zeta function as a part of the wave

function ψ .

The function $\boldsymbol{F(t)}$

The *eta function* is defined for $t \in \mathbb{R}$ and $\sigma > 0$ as the infinite series:

$$\eta(z) = \sum_1^{+\infty} \frac{(-1)^{n-1}}{n^z}$$

or in two infinite series

$$\eta(\sigma + it) = \sum_{n \geq 1} \frac{\cos(t \log n)}{n^\sigma} (-1)^{n-1} - i \sum_{n \geq 1} \frac{\sin(t \log n)}{n^\sigma} (-1)^{n-1}$$

any complex function is also written as

$$\eta(\sigma + it) = \Re \eta(\sigma, t) + i \Im \eta(\sigma, t)$$

The eta function has the same *non trivial* roots of the zeta function ζ , and if we assume the veracity of the **RH**, they are all of the form:

$$\frac{1}{2} + i\gamma_n$$

Waves as a string's vibrations

As the electron is seen as a wave according to the **de Broglie** equation(1924)

$\lambda = \frac{h}{p}$, it is known that the **Schrödinger** equation (1927) , or operator:

$$i\hbar\psi_t = \hat{H}(\psi)$$

$$\psi(X = (x, y, z), t) = f + ig$$

came to model the *wave function* ψ .

As we know from the theory of analytic functions , both $\Re\eta$ and $\Im\eta$ are harmonic functions, verifying the **Laplace** PDE :

$$\frac{\partial^2 \Re\eta(\sigma, t)}{\partial \sigma^2} + \frac{\partial^2 \Re\eta(\sigma, t)}{\partial t^2} = 0$$

Where :

$$\frac{\partial^2 \Re\eta(\sigma, t)}{\partial \sigma^2} = \sum_{n \geq 1} \frac{\cos(t \log n)}{n^\sigma} \times \log^2 n \times (-1)^{n-1}$$

$$\frac{\partial^2 \Re\eta(\sigma, t)}{\partial t^2} = - \sum_{n \geq 1} \frac{\cos(t \log n)}{n^\sigma} \times \log^2 n \times (-1)^{n-1}$$

Nothing prevents us to see $\Delta u = 0$ as a *complex wave equation*:

$$(\Re\eta)_{\sigma\sigma} = i^2(\Re\eta)_{tt}$$

otherwise as

$$i(\Re\eta)_{\sigma\sigma} = -i(\Re\eta)_{tt}$$

For simplicity, we consider here the **Schrödinger** equation for a plane wave (only $X = x$), where \hat{H} is the **Hamiltonian**

$$\hat{H} = -\frac{\hbar^2}{2m}\Delta + V^1$$

We use one dimension wave function $\psi(\mathbf{x}, t)$, and for convenience, we have to interchange the variables in our infinite series $\eta(z)$:

t of $\Re\eta$ is the space x of $\psi : -\infty < t < +\infty$

σ of $\Re\eta$ is the time t of $\psi : 0 < \sigma < +\infty$

the equations gives

¹ "(r, t) designates a potential energy here. For example, it may be the product of an electric potential and the particle's charge. In quantum mechanics, $V(r, t)$ is commonly called a potential. " **Tannoudji – Diu – Laloe** (1/20).

$$\begin{cases} \hbar\psi_\sigma = (\Re\eta)_{\sigma\sigma} & \rightarrow \hbar\psi = (\Re\eta)_\sigma + i\mathbf{F}(t) \\ \hat{H}(\psi) = -i(\Re\eta)_{tt} & \rightarrow \hat{H}\left(\frac{(\Re\eta)_\sigma + iF(t)}{\hbar}\right) = -i(\Re\eta)_{tt} \end{cases}$$

Then the function to be found is

$$\mathbf{F}(t)$$

and we get an expression for the wave function:

$$\psi(t, \sigma) = \frac{(\Re\eta)_\sigma(\sigma, t) + i\mathbf{F}(t)}{\hbar}$$

or more exactly:

$$\psi(\sigma, t) = \frac{i\mathbf{F}(t)}{\hbar} + \frac{1}{\hbar} \sum_{n \geq 1} \frac{\cos(t \log n)}{n^\sigma} \times \log n \times (-1)^{n-1}$$

with , of course, the *normalization* condition:

$$\int_{\mathbb{R}} \|\psi\|^2 dt = 1$$

otherwise :

$$\int_{\mathbb{R}} \left| \sum_{n \geq 1} \frac{(-1)^{n-1}}{n^\sigma} \times \log n \times \cos(t \log n) + i \mathbf{F}(t) \right|^2 dt = \hbar^2$$

If $\int_0^T (\Re \eta)_\sigma^2 dt$ converges at all at $+\infty$. For $(\Re \eta)_\sigma$ we know that:

$$\int_0^T (\Re \eta)_\sigma dt = \sum_1^{+\infty} \frac{(-1)^{n-1}}{n^\sigma} \times \sin(T \log n) < +\infty$$

If there is a problem of convergence, then to avoid it:

- i) We can consider integration in $[0, T]$ for $T > 0$.
- ii) The probability of the existence of a particle diminishes adequately to 0 as $T \rightarrow +\infty$.

Anyway, we get

$$\int_{-T}^T \mathbf{F}(t)^2 dt = \hbar^2 - \int_{-T}^T (\Re \eta)_\sigma^2 dt$$

Since \cos is even then :

$$\int_0^T \mathbf{F}(t)^2 dt = \frac{\hbar^2}{2} - \int_0^T (\Re \eta)_\sigma^2 dt$$

Or otherwise

$$\int_0^T \mathbf{F}(t)^2 dt = \frac{\hbar^2}{2} - \int_0^T (\Re[(1 - 2^{1-z})\zeta(z)])_{\sigma}^2 dt$$

There is some inconsistency here , namely that \hbar^2 is very small. But we can avoid this inconsistency by thinking differently : **we search for possible zeros**, as T grows, of

$$\frac{\hbar^2}{2} - \int_0^T (\Re[(1 - 2^{1-z})\zeta(z)])_{\sigma}^2 dt$$

which , of course, should not be simple zeros (since $\int_0^T \mathbf{F}(t)^2 dt \geq 0$), and ask for their meaning in this situation.

The Hamiltonian in action

Applying the operator on ψ

$$\left(-\frac{\hbar^2}{2m}\Delta + \mathbf{V}\right) \left(\frac{(\Re\eta)_{\sigma} + i\mathbf{F}(t)}{\hbar}\right) = -i(\Re\eta)_{tt}$$

gives the equation

$$-\frac{\hbar^2}{2\hbar m}((\Re\eta)_{\sigma tt} + iF_{tt}(t)) + \frac{1}{\hbar}V((\Re\eta)_{\sigma} + i\mathbf{F}(t)) = -i(\Re\eta)_{tt}$$

where

$$(\Re\eta)_{\sigma tt} = -\sum_{n \geq 1} \frac{\cos(t \log n)}{n^{\sigma}} \times \log^3 n \times (-1)^{n-1}$$

if we want the equation to have sense, then it must verify two conditions :

$$\begin{cases} \frac{\hbar^2}{2m}V((\Re\eta)_{\sigma}) + V(\mathbf{F}(t)) = -\hbar(\Re\eta)_{tt} \\ -\frac{\hbar^2}{2m}(\Re\eta)_{\sigma tt} + V((\Re\eta)_{\sigma}) = 0 \end{cases}$$

The first condition

We know that

$$\eta'(z) = (\Re\eta)_{\sigma} + i(\Im\eta)_{\sigma} = (\Im\eta)_t - i(\Re\eta)_t$$

and

$$\zeta'(z) = \frac{\eta'(z)}{(1 - 2^{1-z})} + \eta(z) \times \left(\frac{1}{1 - 2^{1-z}}\right)'$$

So in the critical strip (let $\sigma = 1$ alone) $\eta'(z)$ and $\zeta'(z)$ have the same zeros. And if **RH** is true then $\eta'(z)$ has no roots in $0 < \sigma < \frac{1}{2}$ because it has been proved that **RH** is equivalent to the fact that $\zeta'(z)$ has no roots when $0 < \sigma < \frac{1}{2}$.

So in the first condition, the term $(\Re\eta)_\sigma$ vanishes at the zeros of $\eta'(z)$ and $\zeta'(z)$. The term $(\Re\eta)_{tt}$ vanishes too for $\sigma = \sigma_0$ and (using **Schwarz**) its zeros are between different vertical sets of zeros. What relation has this first condition with **RH**?

So when does:

$$V(\mathbf{F}(t)) = 0 \quad ?$$

Here V is independent of time. What is the meaning of this in physics?

The second condition

We have seen that the term $(\Re\eta)_\sigma$ vanishes at the zeros of $\eta'(z)$ and $\zeta'(z)$. The term $(\Re\eta)_{\sigma tt}$ vanishes at the zeros of $\eta'(z)$ and $\zeta'(z)$, and this for $\sigma = \sigma_0$ and its zeros are also between different vertical sets of zeros. So the equation

$$V((\Re\eta)_\sigma) = \frac{\hbar^2}{2m} (\Re\eta)_{\sigma tt}$$

Cannot be identically zero.